

Persistent Currents versus Phase Breaking in Mesoscopic Metallic Samples

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- Introduction
- Established theory
- Relation to the dephasing problem?

Eckern, Schwab JLTP **126**, 1291 (2002)

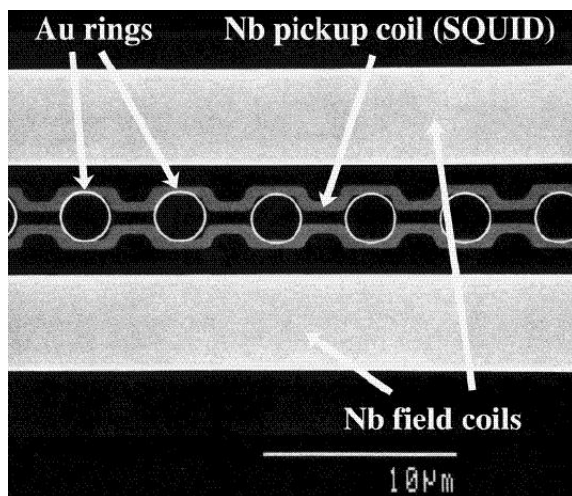
Electron phase coherence in disordered metals

- standard analysis:
electron-phonon interaction, Coulomb interaction
- problem:
 τ_φ often shorter than expected, e.g.
In₂O_{3-x}, heavily doped Si: $1/\tau_\varphi \sim T$
Au (sometimes): $1/\tau_\varphi \sim \sqrt{T}$
many materials: low temperature saturation of τ_φ
- microscopic origin of strong phase breaking?
extreme dilute magnetic impurities?

Persistent currents in normal conducting rings

- phase coherent motion of electrons necessary
—→ sensitive to phase breaking processes
- problem:
sometimes larger than expected theoretically; “wrong” sign

Mohanty, Ann. Phys. '99; Jariwala et al, PRL '01

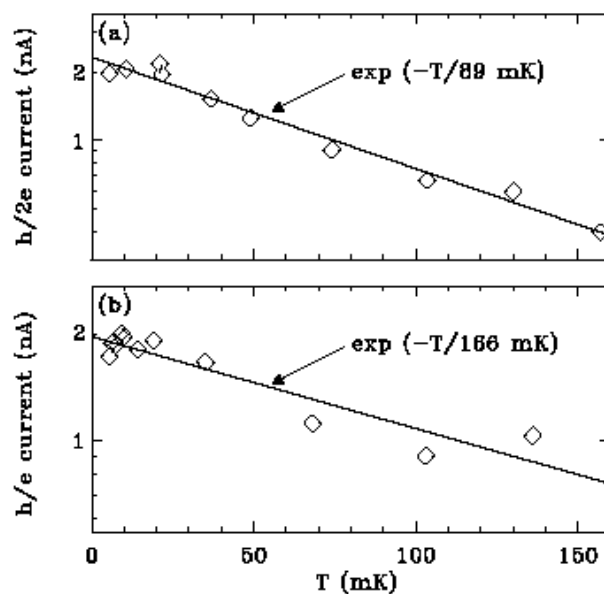
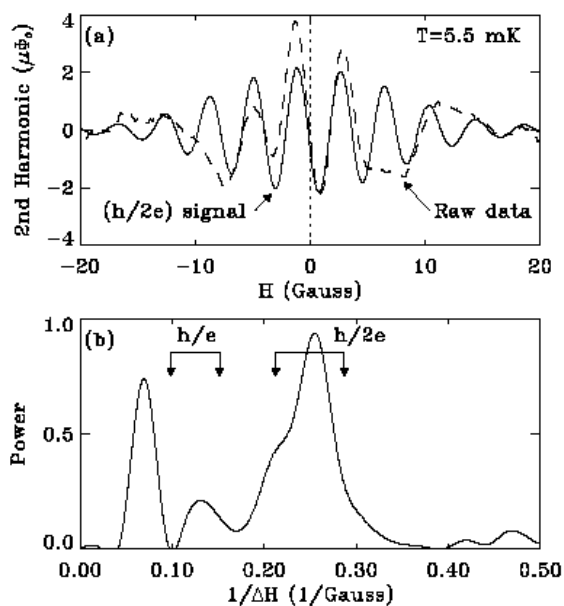


30 gold rings:

$$I \approx e/\tau_d$$

$$\hbar/\tau_d = \hbar v_{F1}/3L^2 \approx 7\text{mK}$$

$$L \approx 8\mu\text{m}, l \approx 87\text{nm}$$



Free fermions

persistent current

$$I = -\frac{\partial \Omega(\phi)}{\partial \phi}, \quad I = -\frac{\partial F(\phi)}{\partial \phi}$$

density of states determines thermodynamics

$$\Omega(\phi) = -2k_B T \int d\epsilon N(\epsilon, \phi) \ln[1 + e^{-(\epsilon - \mu)/k_B T}]$$

density of states for disordered metals

$$N(\epsilon, \phi) = \langle N(\epsilon) \rangle + \delta N(\epsilon, \phi)$$

Fluctuations in the density of states

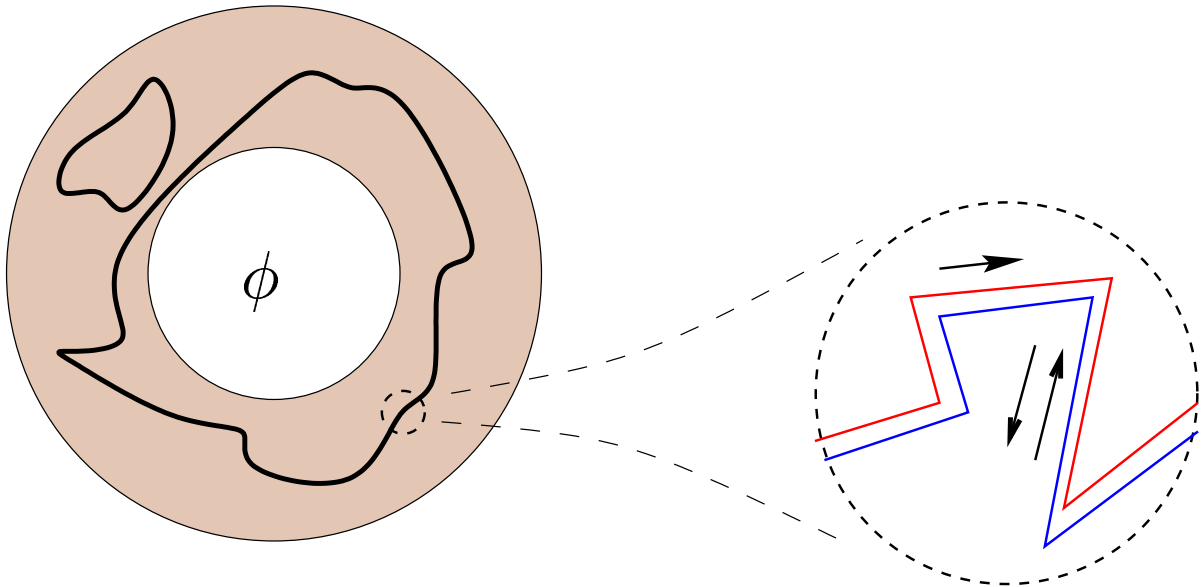
Altshuler & Shklovskii '87

Perturbation theory in $1/\epsilon_F\tau$, $\delta/(\hbar/\tau_d)$ (diffuson, cooperon)

$$\langle N(\epsilon + \omega, \phi) N(\epsilon, \phi') \rangle \propto \sum_q [P(q, \omega)]^2$$

$$\propto \int_0^\infty dt e^{i\omega t} t \left\{ P(0, t) + 2 \sum_m P(mL, t) \cos \left(2\pi m \frac{\phi \pm \phi'}{\phi_0} \right) \right\}$$

$$P(q, \omega) = \frac{1}{-i\omega + Dq^2}$$



$$P(L, t) \propto TT^*$$

magnetic field

$$e^{i(\eta+2\pi\phi/\phi_0)} [e^{i(\eta\pm 2\pi\phi/\phi_0)}]^*$$

Persistent current for free fermions

Cheung et al '89, . . .

average value:

$$\langle I \rangle = -\frac{\partial \langle \Omega \rangle}{\partial \phi} = 0$$

statistical fluctuations:

$$\begin{aligned} \langle \Omega(\phi) \Omega(\phi') \rangle &= \sum_m C_m \cos\left(\frac{2\pi m \phi}{\phi_0}\right) \cos\left(\frac{2\pi m \phi'}{\phi_0}\right) \\ C_m &= \frac{8L}{\pi^2} \int_0^\infty \frac{dt}{t} \left(\frac{\pi k_B T}{\sinh(\pi k_B T t / \hbar)} \right)^2 P(mL, t) \end{aligned}$$

$T = 0$, typical current $I \approx 2\pi \sqrt{C_1} / \phi_0 \approx 3.1e / \tau_d$

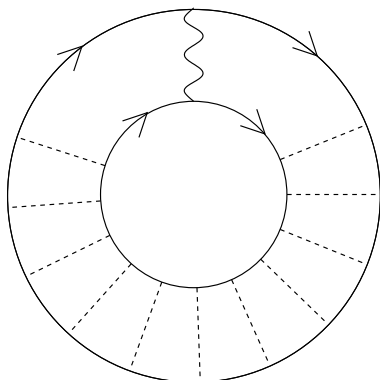
$T \neq 0$: exponential cut-off in long paths

$$\frac{1}{\sinh(\pi k_B T t / \hbar)} \approx 2 \exp(-\pi k_B T t / \hbar)$$

temperature dependence on the scale $k_B T \approx \hbar / \tau_d$

Persistent current, weakly interacting system

Ambegaokar, Eckern '90



$$\delta\Omega_C = \frac{1}{2} \int d\mathbf{r} d\mathbf{r}' v(\mathbf{r}-\mathbf{r}') \delta n(\mathbf{r}) \delta n(\mathbf{r}')$$

$$N_0 \langle v(\mathbf{k}-\mathbf{k}') \rangle_{FS} = \mu_0 \leq 1/2$$

$$\begin{aligned} \langle \delta\Omega_C(\phi) \rangle &= \mu_0 \frac{4L}{\pi\hbar} \int_0^\infty dt \left(\frac{\pi k_B T}{\sinh(\pi k_B T t / \hbar)} \right)^2 \\ &\times \sum_m P(mL, t) \cos(4\pi m \phi / \phi_0) \end{aligned}$$

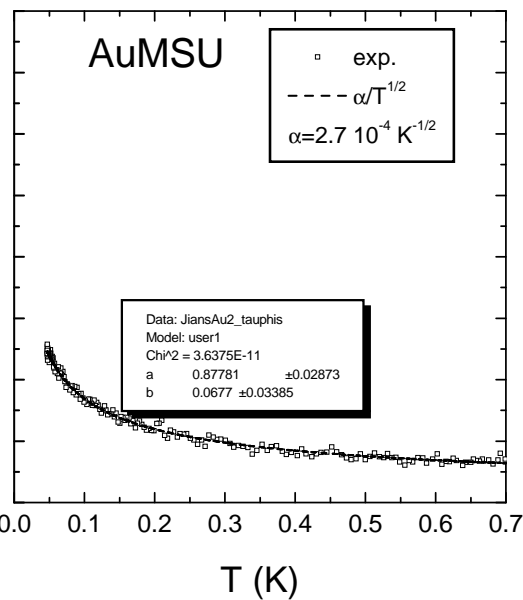
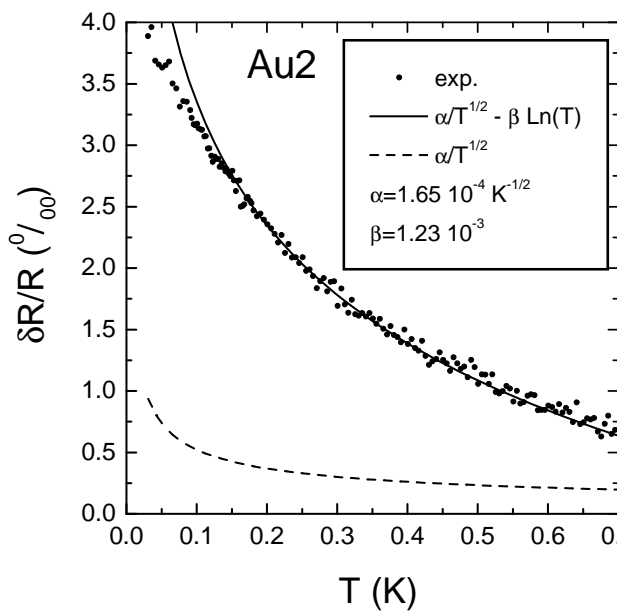
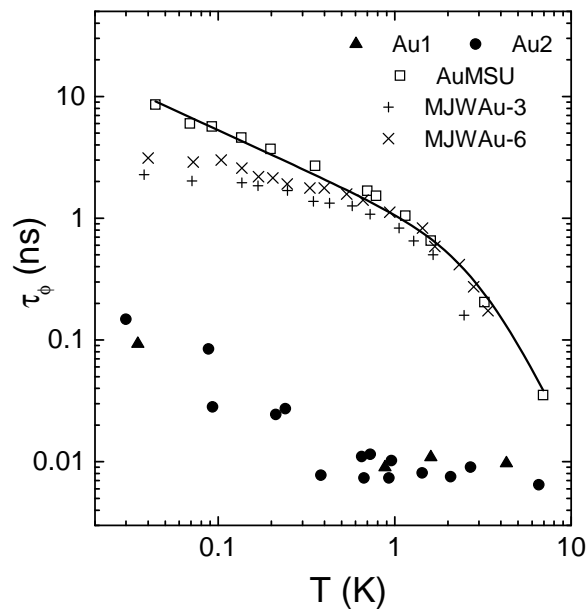
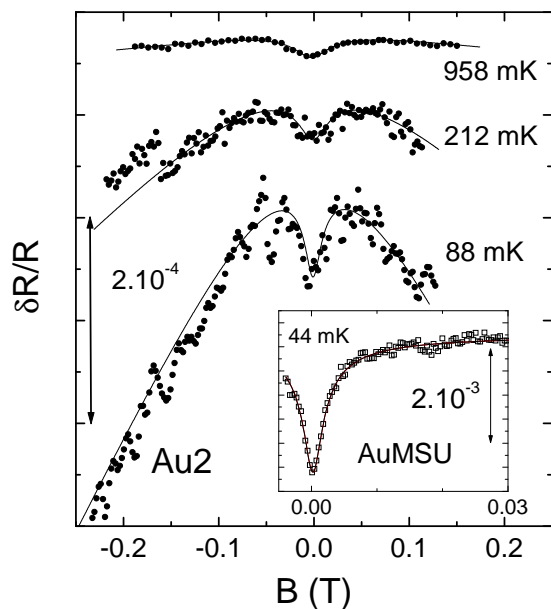
$T = 0$: $I \sim \mu_0 e / \tau_d$, **precise prefactor?**

- exchange term
- renormalized interaction constant
 $\mu_0 \rightarrow \mu^* \approx \mu_0 / \{1 + \mu_0 \ln[\epsilon_F / (\hbar / \tau_d)]\}$
 μ^* depends on the material, $\mu^* < 0.1$ in Cu, Au
- $\mu^* \rightarrow \mu^* - \lambda$?

final result

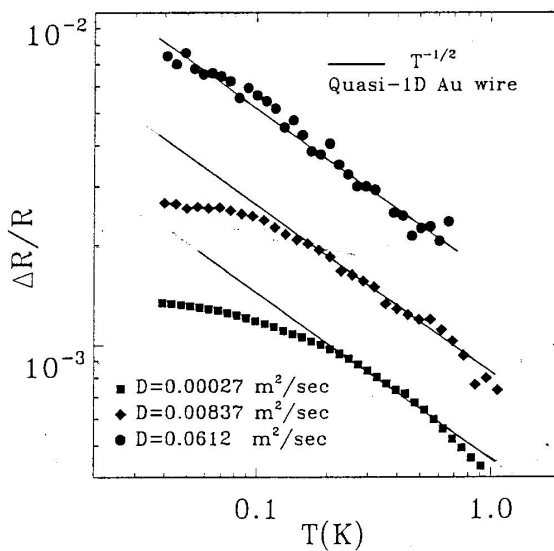
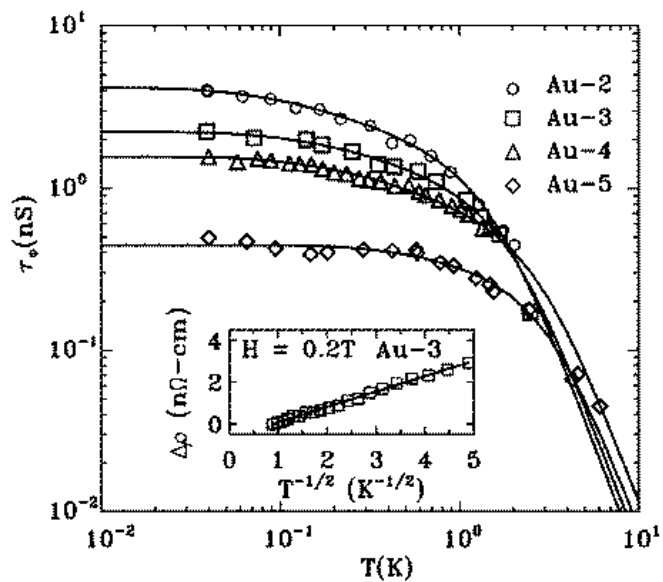
$$I = I_{h/2e} \sin(4\pi\phi/\phi_0) + \dots, \quad I_{h/2e} \approx 8\mu^* / \pi (e/\tau_d)$$

Phase breaking and resistance in gold wires



Pierre, Pothier, Esteve, Devoret, Gougam, Birge;
cond-mat/0012038

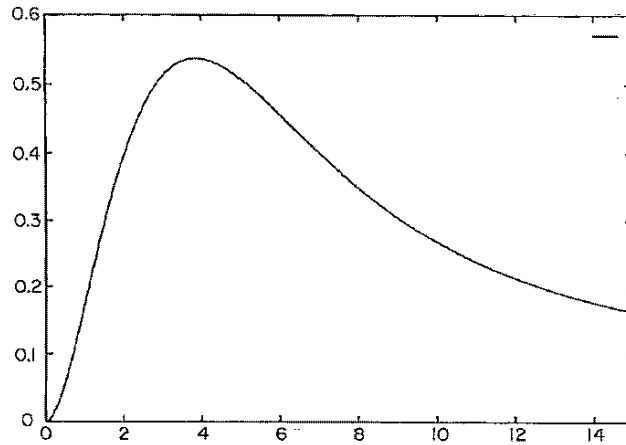
Phase breaking and resistance in gold wires



Mohanty, Jariwala, Webb '97

High frequency electric field

Kravtsov, Yudson '97



$$eE_\omega L \sim \hbar\omega \longrightarrow I_{\text{DC}} \sim (e/\tau_d) \sin(4\pi\phi/\phi_0) + \dots$$

generalization: noise

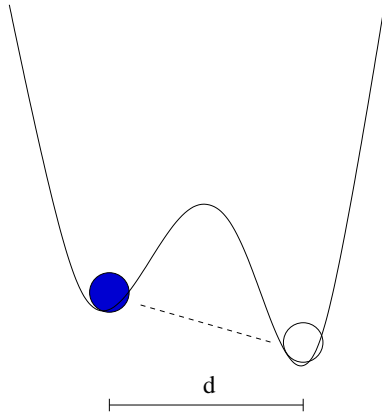
Altshuler, Kravtsov '00

noise: phase breaking & directed current

$$I_{h/2e} = \begin{cases} -(4/\pi)(e/\tau_\phi)e^{-L/L_\phi} & \tau_{\text{so}} \gg \tau_d \\ (2/\pi)(e/\tau_\phi)e^{-L/L_\phi} & \tau_{\text{so}} \ll \tau_d \end{cases}$$

in all experiments: $L \sim L_\phi \longrightarrow I \sim e/\tau_d$.

Electron-impurity interaction?



$$\hat{H}_{\text{TLS}} = \begin{pmatrix} \epsilon & \Delta \\ \Delta & -\epsilon \end{pmatrix}$$

assumption: $P(\epsilon, \Delta) \propto 1/\Delta$

phase breaking:

$$\frac{1}{\tau_\phi} \sim \frac{1}{\tau} \left[1 - \left(\frac{\sin k_F d}{k_F d} \right)^2 \right] \frac{\Delta_{\max}}{\epsilon_{\max}} \frac{1}{\ln(\Delta_{\max}/\Delta_{\min})}$$

($\tau_\phi < \Delta_{\max} < k_B T$; weak coupling)

Imry, Fukuyama, Schwab, EPL 47, 608 (1999)

persistent current:

$$\begin{aligned} I_{h/2e} &\sim -|\mu_{\text{TLS}}| e / \tau_d \\ |\mu_{\text{TLS}}| &\sim \frac{\hbar / \tau_\phi}{\Delta_{\max}} \ln(\Delta_{\max} / \Delta_{\min}) \end{aligned}$$

Schwab, EPJ B 18, 198 (2000)

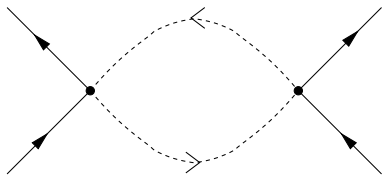
TLS relevant for phase breaking \longrightarrow relevant for persistent current

Persistent current & magnetic defects

Schwab, Eckern '97

$$H = H_0 - \sum_{\text{magnetic defects}} J \mathbf{s}(\mathbf{x}_j) \cdot \mathbf{S}_j$$

effective interaction



$$V_{\text{eff}} \propto \frac{n_s}{N_0^2} \begin{cases} (N_0 J)^2 \chi(T) & T > T_K \\ T_K^{-1} & T < T_K \end{cases}$$

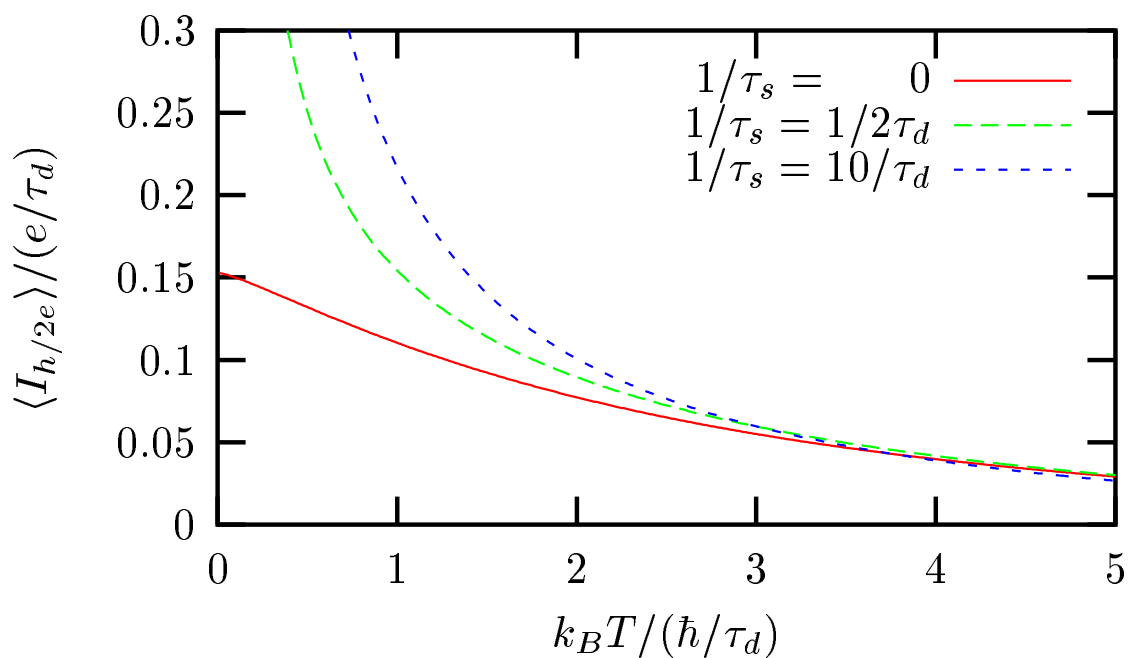
phase breaking

$$\hbar/\tau_s \propto \frac{n_s}{N_0} \begin{cases} (N_0 J)^2 & \text{for } T > T_K \\ (T/T_K)^2 & \text{for } T < T_K \end{cases}$$

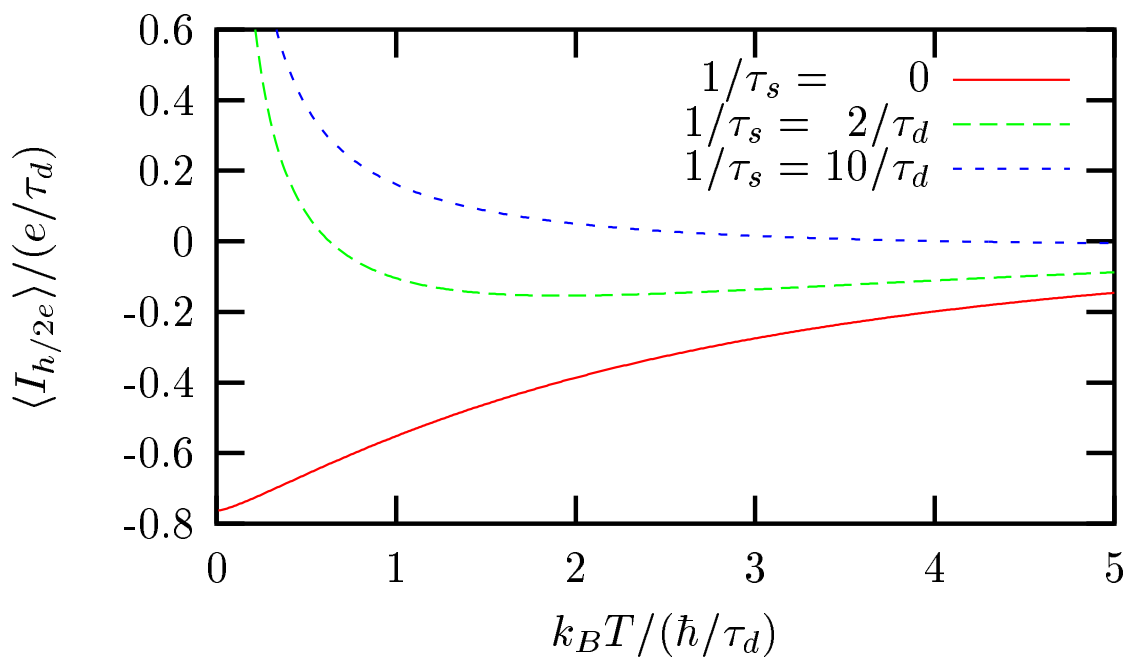
persistent current

$$I \sim N_0 V_{\text{eff}} e / \tau_d$$
$$I \sim \frac{\hbar/\tau_s}{k_B T} e / \tau_d$$

$\mu^* = 0.06$ (Coulomb!) + magnetic defects



$\mu^* = -0.3$ (experiment?) + magnetic defects



Summary

persistent currents in mesoscopic rings are partially understood:

- periodicity $h/e, h/2e$
- characteristic scale e/τ_d
- characteristic temperature scale $k_B T \sim \hbar/\tau_d$
- not understood: sign of $I_{h/2e}$

phase breaking versus persistent currents:

- higher amplitude possible, different T -dependence
- magnetic impurities should be visible