#### The Density Matrix Renormalization Group

- NRG/real space RG
- Particle in a box
  - Problems with NRG
  - Solutions
  - DMRG for I ptle
- Interacting Systems
  - Density Matrix idea
  - DMRG finite system algorithm
  - A few examples from spin chains
- QI perspective: Entanglement and Schmidt Decomposition
- Matrix Product states and Diagrams
  - Periodic BC algorithm





#### DMRG (continued)

- Efficiency
  - Efficient H psi
  - Wavefunction transformation
- Errors, extrapolation
- Two dimensions
  - -TV scan
  - PEPS
  - Finite size scaling with cylindrical BCs





#### Wilson's numerical RG for a Kondo impurity



Standard Feynman diagrammatic perturbation approaches failed in the 60's.

Successes:

- "Poor man's scaling", Anderson et. al. 1970
- •Wilson's NRG, 1975
- •Andrei's exact Bethe ansatz solution, 1982

Wilson's logarithmic basis



#### Wilson's numerical RG



site

Treat short distance, high energy scales first



Diagonalize block, keep m lowest energy states



Add one site, diagonalize block Hamiltonian again, keeping m states



Key point:

Keep track of H through m x m operator and transformation matrices



# Wilson's approach applied in real space



Any truncation yields "kinks" at larger scales.



Wilson's algorithm

- •Diagonalization of entire system
- •Construction of density matrix for block
- •Transformation to new density matrix states
- •Sweeps back and forth





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```
Matrix H(length, length); H = 2.0;
    for(int i = 1: i < n: i++)</pre>
        H(i,i+1) = H(i+1,i) = -1.0;
    Vector leftpsi(1), rightpsi(n-3), evals;
    rightpsi = 1.0; leftpsi = 1.0;
// Finite System sweeps
    for(int it = 1; it <= niter; it++)</pre>
        for(int i = 1; i <= n-3; i++)
            leftpsi = leftpsi / Norm(leftpsi):
            rightpsi = rightpsi / Norm(rightpsi);
            Matrix Htil(4,4), O(n,4); O = 0.0;
            0.Column(1).SubVector(1,i) = leftpsi;
            0(i+1,2) = 0(i+2,3) = 1.0;
            0.Column(4).SubVector(i+3,n) = rightpsi;
            Htil = 0.t() * H * 0;
            EigenValues(Htil,evals,evecs);
            Vector psitil(evecs.Column(1)), newpsi(n);
            newpsi.SubVector(1,i) = leftpsi * psitil(1);
            newpsi(i+1) = psitil(2);
            newpsi(i+2) = psitil(3);
            newpsi.SubVector(i+3,n) = rightpsi * psitil(4);
            if(newpsi(i+1) < 0.0) newpsi *= -1.0;
            cout << "@" << endl;
            for(int j = 1; j <= n; j++)</pre>
                cout << j SP newpsi(j) << endl;</pre>
            if(i < n-3)
                leftpsi = newpsi.SubVector(1,i+1),
                rightpsi = newpsi.SubVector(i+4,n);
            }
```







# S=1 Heisenberg Chain

- Disordered, Haldane gap, finite  $\xi$
- Good picture: Affleck-Kennedy-Lieb-Tasaki state



- AKLT is a matrix product state with m=2 !
- S=I magnon excitations
- S=1/2 free end spins



# Convergence for S=1/2 chain



Comparison with Bethe Ansatz



#### Matrix Product States (Ostlund and Rommer, 1995)

 Insert density matrix/Schmidt eigenstates between all pairs of sites

Matrix Product State:  $\Psi(s_1, s_2, .., s_N) \approx A^1[s_1] A^2[s_2] ... A^N[s_N]$ 

Basic Unit: tensor/matrix

$$A^{s}_{ij} = i - j$$

DMRG Blocks = set of basis states:

$$|s_1\rangle |s_2\rangle |s_3\rangle$$

$$= \{ |i\rangle \}$$



DMRG versus MPS: blocks and bases versus variational states





#### Energy extrapolation



Probability of states thrown away



#### Typical extrapolation of magnetization



Pinning AF fields applied to edges, cylindrical BCs



#### Typical extrapolation of magnetization



High accuracy points indicate quadratic approach!

Pinning AF fields applied to edges, cylindrical BCs



#### Typical extrapolation of magnetization



Pinning AF fields applied to edges, cylindrical BCs



#### Cubic fit to well-converged measurements







- Cylindrical BCs: periodic in y, open in x
- Strong AF pinning fields on left and right edges
- 21 sweeps, up to m=3200 states, 80 hours



Improved finite size scaling: choosing aspect ratios to reduce finite size effects

Μ



Short: proximity to strong pinning makes M large

- "Standard" measurements in QMC estimate  $M^2$  using correlation functions and have large finite size effects  $O(1/L_y)$
- Can one choose a special aspect ratio to eliminate  $O(1/L_y)$  term?
- What is behavior at large length scales? Use finite system spin wave theory as a guide.



#### Square lattice





#### Finite size spin wave theory



- Optimal choice  $\alpha = 1.764$  eliminates linear term
- Even  $\alpha = 1$  has much smaller finite size effects



#### Tilted square lattice



• Tilted lattice has smaller DMRG errors for its width

**1** 0.45

• For this "32x8" obtain M = 0.3052(4)



#### Tilted square lattice



 Results are consistent with and with comparable accuracy to QMC!

#### Results with Sz conservation turned off





- Only one sublattice pinned, other two rotate in a cone
- Other two have z component -M/2
- Here only have  $L_y = 3, 6, 9, ...$





What is the best  $\alpha$  for the triangular lattice? With limited number of widths and finite size effects, could use analytic help...

If  $\alpha$  is within range 1.5 - 2, width 9 data has a range of 0.195 - 0.22, consistent with GFMC and series expansions (but not SWT).

#### Preliminary Results--Triangular lattice

- Current result: M = 0.205(15)
- Consistent with Series, GFMC
- SWT not nearly as accurate as for square lattice

Method	Ref.	N	$E_0/N$	M
Series	this work	$\infty$	-0.5502(4)	0.19(2)
ED	$5,\!68$	12	-0.6103	
		36	-0.5604	0.40
GFQMC	41	$\infty$	-0.5458(1)	0.205(10)
SWT+1/S	42	8	-0.5466	0.2497
SWT+1/S	43	$\infty$		0.266
Coupled cluster	53	$\infty$		0.2134

