## The Density Matrix Renormalization Group

- NRG/real space RG
- Particle in a box
- Problems with NRG
- Solutions
- DMRG for I ptle
- Interacting Systems
- Density Matrix idea
- DMRG finite system algorithm
- A few examples from spin chains
- Ql perspective: Entanglement and Schmidt Decomposition
- Matrix Product states and Diagrams
- Periodic BC algorithm


## DMRG (continued)

- Efficiency
- Efficient H psi
- Wavefunction transformation
- Errors, extrapolation
- Two dimensions
- TV scan
- PEPS
- Finite size scaling with cylindrical BCs


## Wilson's numerical RG for a Kondo impurity



Standard Feynman diagrammatic perturbation approaches failed in the 60 's.

Successes:

- "Poor man's scaling", Anderson et. al. 1970
-Wilson's NRG, 1975
- Andrei's exact Bethe ansatz solution, 1982


Wilson's logarithmic basis

## Wilson's numerical RG



Treat short distance, high energy scales first


Diagonalize block, keep m lowest energy states


Add one site, diagonalize block Hamiltonian again, keeping m states


Key point:
Keep track of H through m x m operator and transformation matrices

## Wilson's approach applied in real space



Wilson's analysis: try it on a particle in a box!


Any truncation yields "kinks" at larger scales.

## DMRG Algorithm



Wilson's algorithm

DMRG sweeps

- Diagonalization of entire system
- Construction of density matrix for block
- Transformation to new density matrix states
- Sweeps back and forth


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```
    Matrix H(length, length); H = 2.0;
    for(int i = 1; i < n; i++)
    H(i,i+1) = H(i+1,i) = -1.0;
    Vector leftpsi(1), rightpsi(n-3), evals;
    rightpsi = 1.0; leftpsi = 1.0;
// Finite System sweeps
    for(int it = 1; it <= niter; it++)
        {
    for(int i = 1; i <= n-3; i++)
        {
        leftpsi = leftpsi / Norm(leftpsi);
        rightpsi = rightpsi / Norm(rightpsi);
        Matrix Htil(4,4), 0(n,4); 0 = 0.0;
        0.Column(1).SubVector(1,i) = leftpsi;
        O(i+1,2) = O(i+2,3) = 1.0;
    0.Column(4).SubVector(i+3,n) = rightpsi;
    Htil = 0.t() * H * 0;
    EigenValues(Htil,evals,evecs);
    Vector psitil(evecs.Column(1)), newpsi(n);
    newpsi.SubVector(1,i) = leftpsi * psitil(1);
    newpsi(i+1) = psitil(2);
    newpsi(i+2) = psitil(3);
    newpsi.SubVector(i+3,n) = rightpsi * psitil(4);
    if(newpsi(i+1) < 0.0) newpsi *= -1.0;
    cout << "@" << endl;
    for(int j = 1; j <= n; j++)
        cout << j SP newpsi(j) << endl;
if(i < n-3)
            leftpsi = newpsi.SubVector(1,i+1),
            rightpsi = newpsi.SubVector(i+4,n);
}
```



## S=I Heisenberg Chain

- Disordered, Haldane gap, finite $\xi$
- Good picture:Affleck-Kennedy-Lieb-Tasaki state

- AKLT is a matrix product state with $\mathrm{m}=2$ !
- $S=1$ magnon excitations
- $S=I / 2$ free end spins


## Convergence for $S=1 / 2$ chain

2000 site $S=1 / 2$ Heisenberg chain



Comparison with Bethe Ansatz

## Matrix Product States (Ostund and Rommer, 1995)

- Insert density matrix/Schmidt eigenstates between all pairs of sites
Matrix Product State: $\quad \Psi\left(\mathrm{s}_{1}, \mathrm{~s}_{2}, . . \mathrm{s}_{\mathrm{N}}\right) \approx \mathrm{A}^{1}\left[\mathrm{~s}_{1}\right] \mathrm{A}^{2}\left[\mathrm{~s}_{2}\right] \ldots \mathrm{A}^{\mathrm{N}}\left[\mathrm{s}_{\mathrm{N}}\right]$
$2^{\mathrm{N}}$

$\mathrm{N} \mathrm{m}^{2}$

Basic Unit: tensor/matrix

$$
A^{s}{ }_{i j}=\overbrace{i-(A)}^{s}
$$

DMRG Blocks $=$ set of basis states:


DMRG versus MPS: blocks and bases versus variational states
DMRG wavefunction:

New block:


Operators:


## Energy extrapolation



Probability of states thrown away

## Typical extrapolation of magnetization



Pinning AF fields applied to edges, cylindrical BCs

## Typical extrapolation of magnetization



High accuracy points indicate quadratic approach!

Pinning AF fields applied to edges, cylindrical BCs

## Typical extrapolation of magnetization




Pinning AF fields applied to edges, cylindrical BCs

Cubic fit to well-converged measurements


## Square lattice: benchmark against QMC


$20 \times 10$

- Cylindrical BCs: periodic in $y$, open in $x$
- Strong AF pinning fields on left and right edges
- 21 sweeps, up to $m=3200$ states, 80 hours

Improved finite size scaling: choosing aspect ratios to reduce finite size effects


Long: 1D makes M small


Short: proximity to strong pinning makes M large

- "Standard" measurements in QMC estimate $\mathrm{M}^{2}$ using correlation functions and have large finite size effects $O\left(1 / L_{y}\right)$
- Can one choose a special aspect ratio to eliminate $O\left(1 / L_{y}\right)$ term?
- What is behavior at large length scales? Use finite system spin wave theory as a guide.


## Square lattice



## Finite size spin wave theory



- Optimal choice $\alpha=1.764$ eliminates linear term
- Even $\alpha=1$ has much smaller finite size effects


## Tilted square lattice


$\uparrow 0.45$

- Tilted lattice has smaller DMRG errors for its width
- For this " $32 \times 8$ " obtain $M=0.3052(4)$


## Tilted square lattice



Sandvik QMC

- Results are consistent with and with comparable accuracy to QMC!

Results with Sz conservation turned off


## Triangular Lattice



- Only one sublattice pinned, other two rotate in a cone
- Other two have z component -M/2
- Here only have $L_{y}=3,6,9, \ldots$


## Triangular lattice, Scaled Data



What is the best $\alpha$ for the triangular lattice? With limited number of widths and finite size effects, could use analytic help...

If $\alpha$ is within range 1.5-2, width 9 data has a range of $0.195-0.22$, consistent with GFMC and series expansions (but not SWT).

## Preliminary Results--Triangular lattice

- Current result: $M=0.205$ (I5)
- Consistent with Series, GFMC
- SWT not nearly as accurate as for square lattice

| Method | Ref. | $N$ | $E_{0} / N$ | $M$ |
| :--- | :--- | :--- | :--- | :--- |
| Series | this work | $\infty$ | $-0.5502(4)$ | $0.19(2)$ |
| ED | 5,68 | 12 | -0.6103 |  |
|  |  | 36 | -0.5604 | 0.40 |
| GFQMC | 41 | $\infty$ | $-0.5458(1)$ | $0.205(10)$ |
| SWT+1/S | 42 | $\infty$ | -0.5466 | 0.2497 |
| SWT $+1 / \mathrm{S}$ | 43 | $\infty$ |  | 0.266 |
| Coupled cluster | 53 | $\infty$ |  | 0.2134 |

